

# Balancing Domain Decomposition for Mortar Mixed Finite Element Methods

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Solution Methods for Large-Scale Nonlinear Problems  
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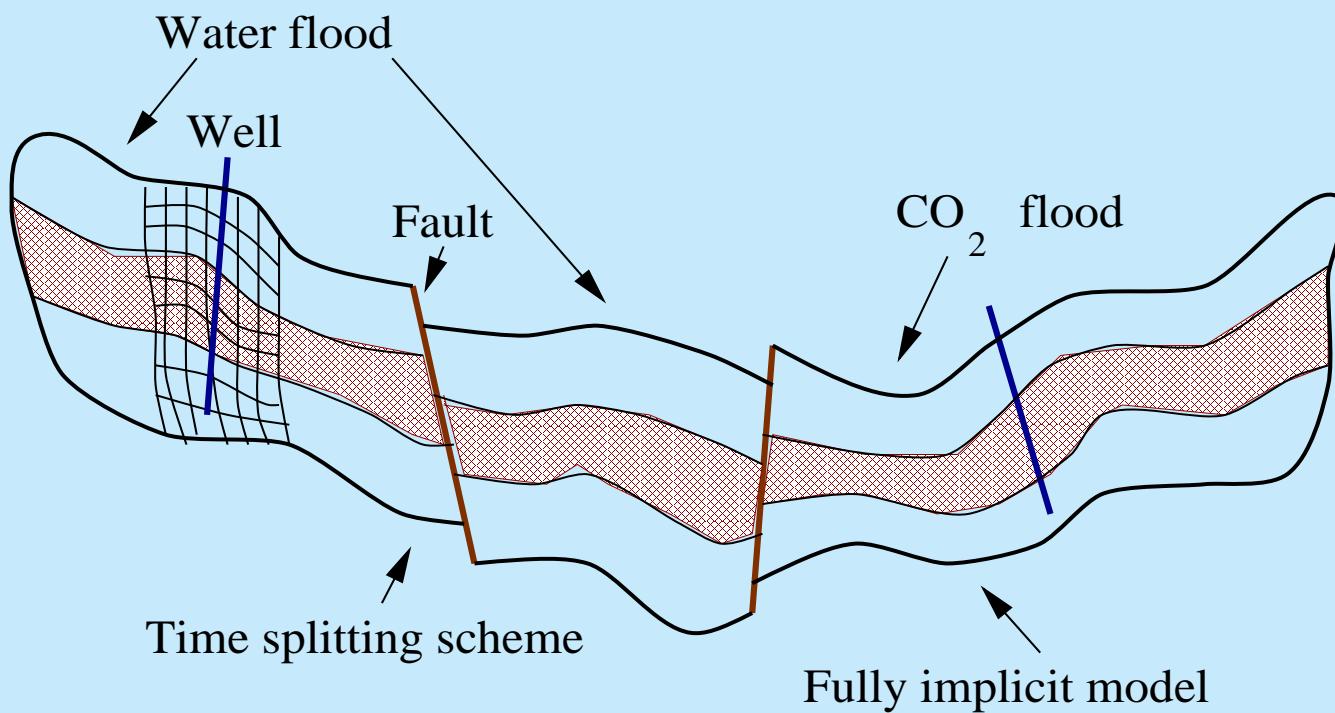
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## Outline

- Multiblock formulation and discretization for single-phase flow
- Non-overlapping domain decomposition
- Balancing preconditioner for the interface problem
- Condition number estimates
- Numerical experiments
- Extensions to multiphase flow
- Interface Newton-GMRES solver

## Multiblock multiphysics approach to flow in porous media



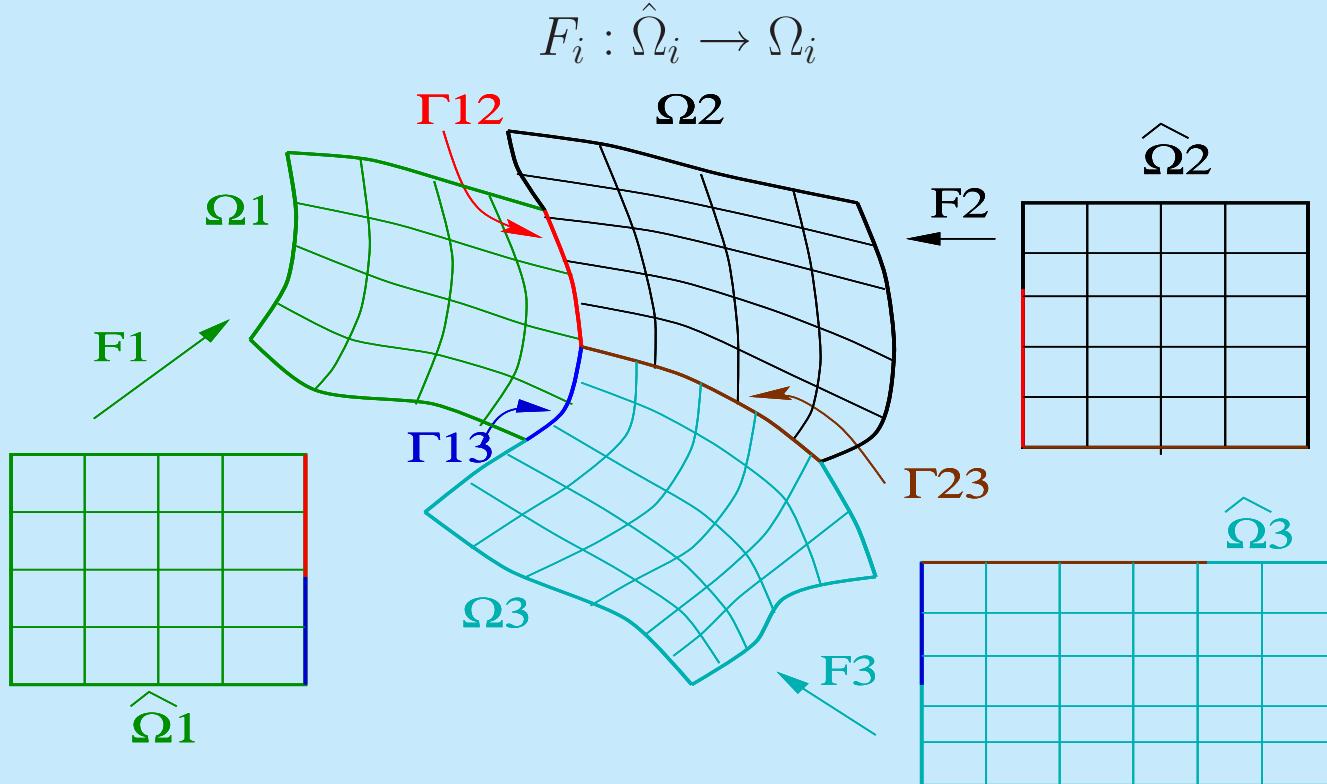
## Multiblock gridding strategy

Partition the domain into a series of blocks:

$$\bar{\Omega} = \bigcup_{i=1}^n \bar{\Omega}_i, \quad \Gamma_{ij} = \partial\Omega_i \cap \partial\Omega_j$$

Each block is covered by a logically rectangular grid.

Consider each block (grid) as the image of a rectangular reference block (grid):



## Features of multiblock subsurface modeling

- Accuracy and local mass conservation
- Irregular geometry - flexible (local) gridding
- Geological features - faults, fractures, layers - non-matching grids
- Multiphysics and local phenomena
- Efficiency - large scale simulations
  - Parallel scalability - domain decompositions
  - Reduced computational cost - different numerical methods on different blocks
  - Local and interface adaptivity

### Critical issues:

- Imposing matching conditions across interfaces in a stable and accurate way
- Developing efficient nonlinear interface solvers and preconditioners

## Single phase flow model

$\Omega \subset \mathbf{R}^d$  ( $d = 2$  or  $3$ ) - flow domain

$$\mathbf{u} = -K\nabla p \quad \text{in } \Omega \quad (\text{Darcy's law})$$

$$\nabla \cdot \mathbf{u} = q \quad \text{in } \Omega \quad (\text{conservation of mass})$$

$$\mathbf{u} \cdot \nu = 0 \quad \text{on } \partial\Omega \quad (\text{no flow BC})$$

$\mathbf{u}$  - velocity,  $p$  - pressure,  $K$  - permeability tensor,  $q$  - source (wells)

## Variational mixed formulation

$$H(\text{div}; \Omega) = \{\mathbf{v} : \mathbf{v} \in (L^2(\Omega))^d, \nabla \cdot \mathbf{v} \in L^2(\Omega)\}$$

$$\mathbf{V} = \{\mathbf{v} \in H(\text{div}; \Omega) : \mathbf{v} \cdot \nu = 0 \text{ on } \partial\Omega\}, \quad W = L^2(\Omega)$$

Find  $\mathbf{u} \in \mathbf{V}$ ,  $p \in W$  such that

$$(K^{-1}\mathbf{u}, \mathbf{v}) = (p, \nabla \cdot \mathbf{v}), \quad \mathbf{v} \in \mathbf{V},$$

$$(\nabla \cdot \mathbf{u}, w) = (q, w), \quad w \in W.$$

## The mixed finite element method

$T_h$  - finite element partition

$\mathbf{V}_h \times W_h \subset \mathbf{V} \times W$  - mixed finite element spaces

Find  $\mathbf{u}_h \in \mathbf{V}_h$ ,  $p_h \in W_h$  such that

$$\begin{aligned}(K^{-1}\mathbf{u}_h, \mathbf{v}) &= (p_h, \nabla \cdot \mathbf{v}), & \mathbf{v} \in \mathbf{V}_h, \\ (\nabla \cdot \mathbf{u}_h, w) &= (q, w), & w \in W_h.\end{aligned}$$

### Properties:

- Simultaneous (accurate) approximation of pressure and velocity
- Local mass conservation: for each element  $E$ ,

$$w = \begin{cases} 1 \text{ on } E, \\ 0 \text{ otherwise} \end{cases} \implies \int_E \nabla \cdot \mathbf{u}_h = \int_E q.$$

- Continuity of normal flux across element faces: for each  $e = \partial E_1 \cap \partial E_2$ ,

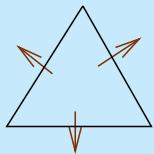
$$\int_e (\mathbf{u}_h|_{E_1} \cdot \nu - \mathbf{u}_h|_{E_2} \cdot \nu) = 0.$$

## Raviart-Thomas spaces

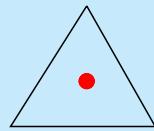
- Triangular element:

$$\mathbf{V}_h^k(E) = (P_k(E))^2 + \mathbf{x}P_k(E)$$

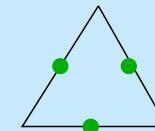
$$W_h^k(E) = P_k(E)$$



↑ velocity



● pressure

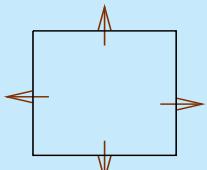


● Lagrange  
multiplier

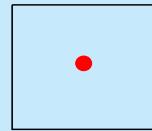
- Rectangular element:

$$\mathbf{V}_h^k(E) = P_{k+1,k}(E) \times P_{k,k+1}(E)$$

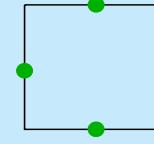
$$W_h^k(E) = P_{k,k}(E),$$



↑ velocity



● pressure



● Lagrange  
multiplier

## Multiblock formulation for single phase flow

$$\bar{\Omega} = \cup_{i=1}^n \bar{\Omega}_i$$

$$\Gamma_{ij} = \partial\Omega_i \cap \partial\Omega_j$$

On each block  $\Omega_i$ :

$$\begin{aligned}\mathbf{u} &= -K\nabla p && \text{in } \Omega_i \\ \nabla \cdot \mathbf{u} &= q && \text{in } \Omega_i \\ \mathbf{u} \cdot \nu &= 0 && \text{on } \partial\Omega_i \cap \partial\Omega\end{aligned}$$

On each interface  $\Gamma_{ij}$ :

$$\begin{aligned}p_i &= p_j && \text{on } \Gamma_{ij} \\ [\mathbf{u} \cdot \nu]_{ij} &= 0 && \text{on } \Gamma_{ij}\end{aligned}$$

where

$$\begin{aligned}p_i &= p|_{\partial\Omega_i} \\ [\mathbf{u} \cdot \nu]_{ij} &\equiv \mathbf{u}|_{\Omega_i} \cdot \nu - \mathbf{u}|_{\Omega_j} \cdot \nu\end{aligned}$$

## Multiblock (macro-hybrid) discretization

**On each block  $\Omega_i$ :**

$\mathcal{T}_{h,i}$  - finite element partition of  $\Omega_i$

$\mathcal{T}_{h,i}$  and  $\mathcal{T}_{h,j}$  - possibly non-matching on  $\Gamma_{ij}$

$(\mathbf{u}_{h,i}, p_{h,i}) \in \mathbf{V}_{h,i} \times W_{h,i}$  - mixed finite element approximation on  $\mathcal{T}_{h,i}$

**On each interface  $\Gamma_{ij}$ :**

Interface finite element grid  $\mathcal{T}_{h,ij}$  - possibly different from the grids on  $\Omega_i$  and  $\Omega_j$ .

Mortar finite element space  $M_{h,ij}$  on  $\mathcal{T}_{h,ij}$ .

Interface pressure is approximated by  $\lambda_h$  in the mortar space  $M_{h,ij}$ .

Flux continuity is imposed weakly:

$$\int_{\Gamma_{ij}} [\mathbf{u} \cdot \nu]_{ij} \mu \, d\sigma = 0, \quad \forall \mu \in M_{h,ij}.$$

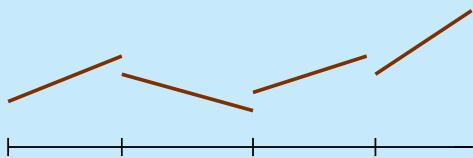
## Mortar finite element spaces for $\mathbf{RT}_0$

**Matching interface  $\Gamma_{ij}$ :** Use the standard piecewise constant Lagrange multiplier space.

**Non-matching interface  $\Gamma_{ij}$ :** The piecewise constant space leads to only sub-optimally convergent method.

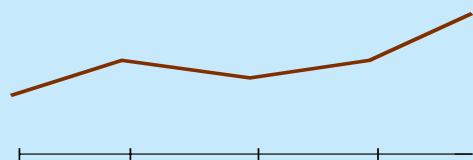
Instead, use piecewise linears:

$$M_{h,ij}^d = \{\mu \in L^2(\Gamma_{ij}) : \mu|_e \text{ is (bi)linear}\}$$



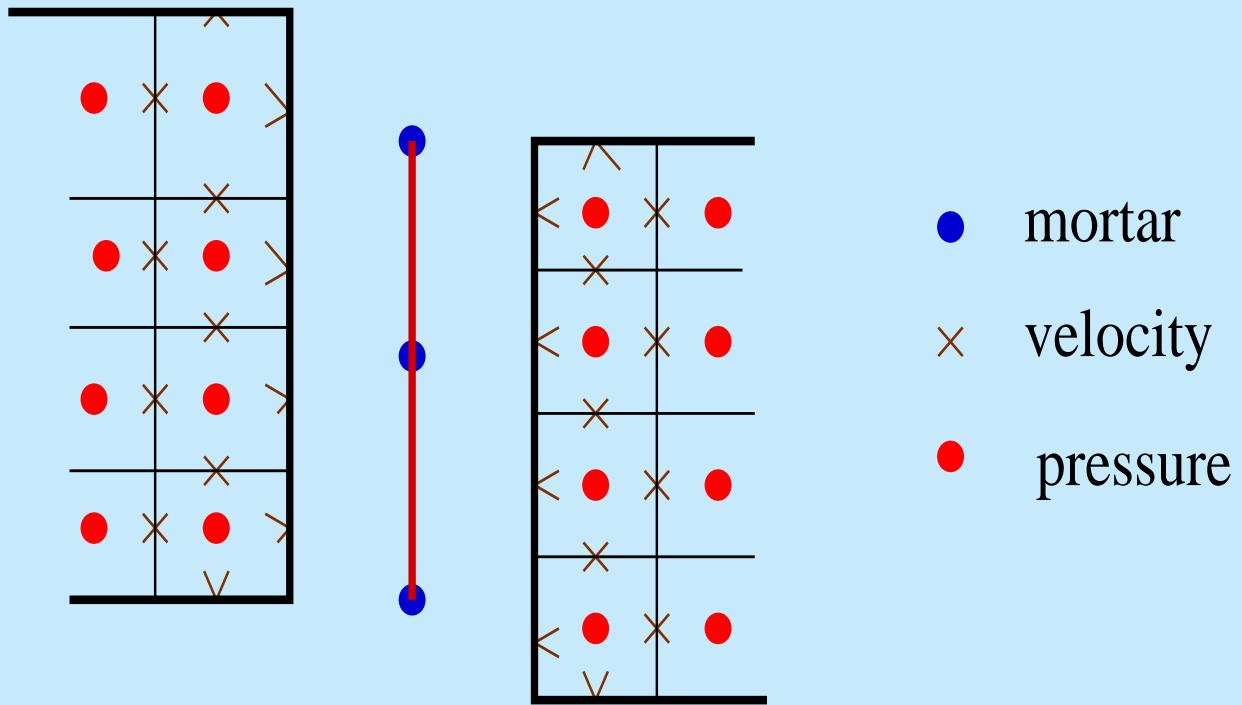
and

$$M_{h,ij}^c = M_{h,ij}^d \cap C^0(\Gamma_{ij}).$$



$M_{h,ij} = M_{h,ij}^d$  or  $M_{h,ij}^c$  - a mortar finite element space on  $\Gamma_{ij}$ .

## Multiblock discretization spaces



$$\mathbf{V}_h = \bigoplus_{i=1}^n \mathbf{V}_{h,i}, \quad W_h = \bigoplus_{i=1}^n W_{h,i}, \quad M_h = \bigoplus_{0 \leq i < j \leq n} M_{h,ij}$$

## The mortar mixed finite element method

Find  $\mathbf{u}_h \in \mathbf{V}_h$ ,  $p_h \in W_h$ ,  $\lambda_h \in M_h$  s.t. for  $1 \leq i \leq n$

$$\begin{aligned} (K^{-1}\mathbf{u}_h, \mathbf{v})_{\Omega_i} &= (p_h, \nabla \cdot \mathbf{v})_{\Omega_i} - \langle \lambda_h, \mathbf{v} \cdot \nu_i \rangle_{\Gamma_i}, & \mathbf{v} \in \mathbf{V}_{h,i}, \\ (\nabla \cdot \mathbf{u}_h, w)_{\Omega_i} &= (q, w)_{\Omega_i}, & w \in W_{h,i}, \\ \sum_{i=1}^n \langle \mathbf{u}_h \cdot \nu_i, \mu \rangle_{\Gamma_i} &= 0, & \mu \in M_h. \end{aligned}$$

**Stability condition on  $M_{h,ij}$ :**

$$\mathcal{Q}_{h,i}\mu = \mathcal{Q}_{h,j}\mu = 0 \Rightarrow \mu = 0, \quad \forall \mu \in M_{h,ij},$$

where  $\mathcal{Q}_{h,i} : M_{h,ij} \rightarrow \mathbf{V}_{h,i} \cdot \nu|_{\Gamma_{ij}}$  is the  $L^2$ -projection:

$$\langle \mu - \mathcal{Q}_{h,i}\mu, \mathbf{v} \cdot \nu \rangle = 0, \quad \mathbf{v} \in \mathbf{V}_{h,i}.$$

**Lemma:** There exists a unique solution  $(\mathbf{u}_h, p_h, \lambda_h)$  such that

$$\|\mathbf{u}_h\|_{\mathbf{V}} + \|p_h\|_W + \|\lambda_h\|_{L^2(\Gamma)} \leq C\|q\|_{L^2(\Omega)}.$$

## Convergence results

**Theorem** (Arbogast, Wheeler, Cowsar, Y.)

Assume that the grids are quasi-uniform and for each  $\mu \in \Lambda_h$

$$\|\mu\|_{0,\Gamma_{i,j}} \leq C(\|\mathcal{Q}_{h,i}\mu\|_{0,\Gamma_{i,j}} + \|\mathcal{Q}_{h,j}\mu\|_{0,\Gamma_{i,j}}).$$

Then

$$\|\mathbf{u} - \mathbf{u}_h\|_0 \leq C \sum_{i=1}^n (\|p\|_{r+1,\Omega_i} + \|\mathbf{u}\|_{r,\Omega_i}) h^r, \quad 0 < r \leq k+1,$$

$$\|\nabla \cdot (\mathbf{u} - \mathbf{u}_h)\|_0 \leq C \sum_{i=1}^n \|\nabla \cdot \mathbf{u}\|_{r,\Omega_i} h^r, \quad 0 < r \leq l+1,$$

$$|||\mathbf{u} - \mathbf{u}_h||| \leq C \sum_{i=1}^n (\|p\|_{r+3/2,\Omega_i} + \|\mathbf{u}\|_{r+1/2,\Omega_i}) h^{r+1/2}, \quad 0 < r \leq k+1,$$

$$\|\hat{p} - p_h\|_0 \leq C \sum_{i=1}^n (\|p\|_{r+1,\Omega_i} + \|\mathbf{u}\|_{r,\Omega_i} + \|\nabla \cdot \mathbf{u}\|_{r,\Omega_i}) h^{r+1},$$

$$\|p - p_h\|_0 \leq C \sum_{i=1}^n (\|p\|_{r+1,\Omega_i} + \|\mathbf{u}\|_{r,\Omega_i} + \|\nabla \cdot \mathbf{u}\|_{r,\Omega_i}) h^r,$$

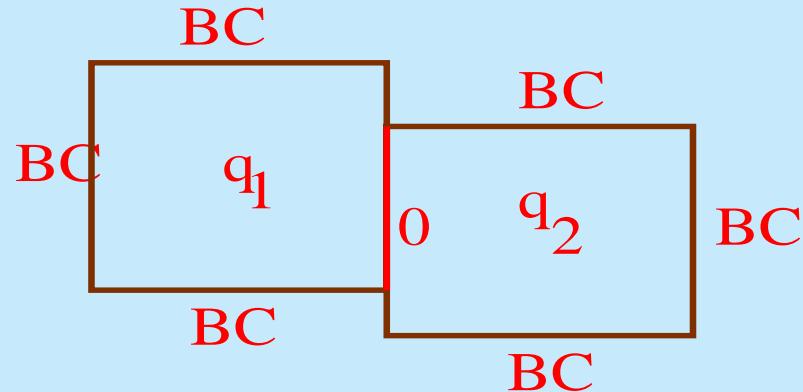
$$0 < r \leq \min(k+1, l+1).$$

Here  $|||\cdot|||$  denote the discrete  $L^2$ -norm along the Gaussian lines.

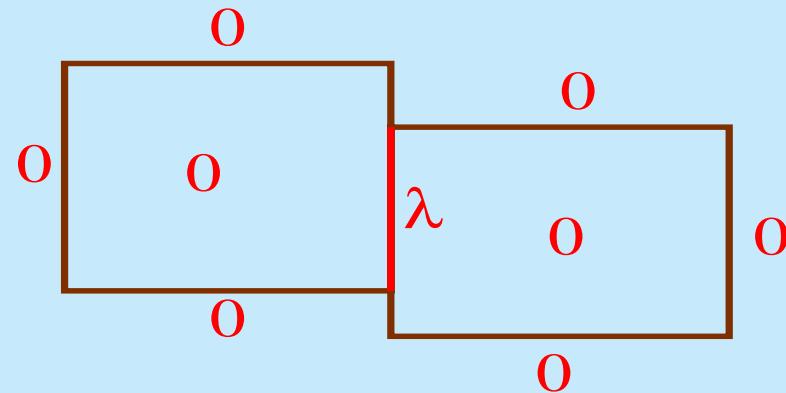
## Parallel domain decomposition

[Glowinski, Wheeler; Y] Two types of subdomain problems:

$(\bar{\mathbf{u}}_h|_{\Omega_i}, \bar{p}_h|_{\Omega_i})$  is the solution on  $\Omega_i$  with zero mortar data, correct source and BC's.



$(\mathbf{u}_h^*(\lambda)|_{\Omega_i}, p_h^*(\lambda)|_{\Omega_i})$  is the solution on  $\Omega_i$  with mortar data  $\lambda \in M_h$ , zero source and BC's.



## Reduction to an interface problem

Define  $s_h : L^2(\Gamma) \times L^2(\Gamma) \rightarrow \mathbf{R}$  by

$$s_h(\lambda, \mu) = -\sum_{i,j} \int_{\Gamma_{i,j}} [\mathbf{u}_h^*(\lambda) \cdot \nu]_{ij} \mu d\sigma$$

Define  $g_h : L^2(\Gamma) \rightarrow \mathbf{R}$  by

$$g_h(\mu) = \sum_{i,j} \int_{\Gamma_{i,j}} [\bar{\mathbf{u}}_h \cdot \nu]_{ij} \mu d\sigma$$

The solution  $(\mathbf{u}_h, p_h, \lambda_h)$  to the original problem satisfies

$$s_h(\lambda_h, \mu) = g_h(\mu), \quad \mu \in M_h,$$

with

$$\mathbf{u}_h = \mathbf{u}_h^*(\lambda_h) + \bar{\mathbf{u}}_h, \quad p_h = p_h^*(\lambda_h) + \bar{p}_h.$$

### **Lemma**

$s_h(\cdot, \cdot)$  is symmetric and positive semi-definite on  $M_h$ .

Solve the interface problem

$$s_h(\lambda_h, \mu) = g_h(\mu), \quad \mu \in M_h$$

in parallel by using conjugate gradient or multigrid.

## Iterative interface solver

Evaluation of the bilinear form on each iteration:

- Project  $\lambda$  onto  $\mathbf{V}_{h,i} \cdot \nu$ :  $\lambda \rightarrow \mathcal{Q}_{h,i}\lambda$
- Solve in parallel subdomain problems with Dirichlet data  $\mathcal{Q}_{h,i}\lambda$  on the interior interfaces to compute the fluxes  $\mathbf{u}_h^*(\lambda) \cdot \nu$
- Project the fluxes back to the mortar space  $M_h$  to compute the jumps across the interfaces

## Algebraic system

$$\begin{pmatrix} A & B & L \\ B^T & 0 & 0 \\ L^T & 0 & 0 \end{pmatrix} \begin{pmatrix} \mathbf{u} \\ p \\ \lambda \end{pmatrix} = \begin{pmatrix} 0 \\ q \\ 0 \end{pmatrix} \Leftrightarrow \begin{pmatrix} R & L \\ L^T & 0 \end{pmatrix} \begin{pmatrix} \mathbf{x} \\ \lambda \end{pmatrix} = \begin{pmatrix} \tilde{q} \\ 0 \end{pmatrix}$$

$\mathbf{x}$  - subdomain unknowns;  $\lambda$  - interface unknowns

Matrix is symmetric but indefinite. Form the Shur complement system:

$$L^T R^{-1} L \lambda = L^T R^{-1} \tilde{q} \quad (1)$$

## Interface problem

Define, for  $i = 1, \dots, n$ ,  $S_i : M_h \rightarrow M_h$  and  $s_i : L^2(\Gamma) \times L^2(\Gamma) \rightarrow \mathbf{R}$  such that for all  $\lambda, \mu \in M_h$

$$\langle S_i \lambda, \mu \rangle = s_i(\lambda, \mu) = -\langle \mathbf{u}_i^*(\lambda) \cdot \nu_i, \mu \rangle_{\Gamma_i}.$$

$S_i$  is a Dirichlet-to-Neumann map:

$$S_i \lambda = -\mathcal{Q}_i^T(\mathbf{u}_i^*(\lambda) \cdot \nu_i)$$

Define  $S : M_h \rightarrow M_h$  and  $s : L^2(\Gamma) \times L^2(\Gamma) \rightarrow \mathbf{R}$  such that

$$S = \sum_{i=1}^n S_i, \quad s(\lambda, \mu) = \sum_{i=1}^n s_i(\lambda, \mu).$$

The interface problem is

$$S\lambda = g \quad \text{or} \quad s(\lambda, \mu) = g(\mu), \quad \forall \mu \in M_h.$$

## Balancing preconditioner

### Balancing preconditioner background

Interface formulation: Glowinski and Wheeler (1988)

Neumann-Neumann preconditioner:

Bourgat, Glowinski, Le Tallec, and Vidrascu (1989)

De Roeck and Le Tallec (1991)

Le Tallec (1993)

Achdou, Le Tallec, Nataf, and Vidrascu (2000)

Balancing preconditioner:

Mandel (1993)

Mandel and Brezina (1996)

Cowsar, Mandel, and Wheeler (1995)

## Interface Neumann-Neumann preconditioner

Define a partition of unity  $D_i$ :  $\text{supp}(D_i) \subset \Gamma_i$  and  $\sum_{i=1}^n D_i \lambda = \lambda, \forall \lambda \in M_h$ .

**ALGORITHM:**

Given  $r \in M_h$ , define  $M_{NN}^{-1}r$  as follows:

1.  $r_i = D_i^T r$  (distribute  $r$  to subdomains).
2. Solve  $S_i \lambda_i = r_i$  for  $\lambda_i \in M_h$ .
3.  $M_{NN}^{-1}r = \sum_{i=1}^n D_i \lambda_i$  (average local solutions).

$$M_{NN}^{-1} = \sum_{i=1}^n D_i S_i^{-1} D_i^T$$

$S_i^{-1} : r_i \rightarrow \lambda_i$  is a Neumann-to-Dirichlet map and requires a local Neumann solve:

Find  $\mathbf{u}_i \in V_{h,i}^0$ ,  $p_i \in W_{h,i}$ ,  $\lambda_i \in M_{h,i}$  such that

$$\begin{aligned} (K^{-1} \mathbf{u}_i, \mathbf{v})_{\Omega_i} &= (p_i, \nabla \cdot \mathbf{v})_{\Omega_i} - \langle \lambda_i, \mathbf{v} \cdot \nu_i \rangle_{\Gamma_i}, & \mathbf{v} \in \mathbf{V}_{h,i}^0, \\ (\nabla \cdot \mathbf{u}_i, w)_{\Omega_i} &= 0, & w \in W_{h,i}, \\ \langle \mathbf{u}_i \cdot \nu_i, \mu \rangle_{\Gamma_i} &= \langle r_i, \mu \rangle_{\Gamma_i}, & \mu \in M_{h,i}. \end{aligned}$$

Problems:  $S_i$  may be singular; No global exchange of information.

## Balancing preconditioner

1. Balance residuals so local problems  $S_i \lambda_i = r_i$  are solvable (modulo Null  $S_i$ ).
2. Result should not depend on the specific choice of local solutions.
3. Parallel scalability - global transfer of information through a coarse solve.

$S_i \lambda_i = r_i$  is solvable if  $r_i \perp \text{Null}S_i$ .

$$\text{Null}S_i = \begin{cases} \{\text{const}\} & \text{if full Neumann} \\ \emptyset & \text{otherwise} \end{cases}$$

Define  $Z_i$ :  $\text{Null } S_i \subseteq Z_i$  (Take  $Z_i = \{\text{const}\}$ ).

Coarse space:

$$M_H = \{\lambda \in M_h : \lambda = \sum_{i=1}^n D_i \zeta_i, \ z_i \in Z_i\}.$$

$\dim M_H \leq n$  (number of subdomains)

A residual  $r$  is balanced (local problems are solvable) if

$$\langle r, \mu_H \rangle_\Gamma = 0, \quad \forall \mu_H \in M_H.$$

Balancing  $r$ :  $r^{bal} = r - Sr_H$  where  $r_H \in M_H$  is found by solving a coarse problem

$$s(r_H, \mu_H) = \langle r, \mu_H \rangle_\Gamma, \quad \forall \mu_H \in M_H.$$

## Balancing preconditioner (Cont.)

ALGORITHM:

Given  $r \in M_h$ , define  $M_{bal}^{-1}r$  as follows:

1. Solve a coarse problem

$$s(r_H, \mu_H) = \langle r, \mu_H \rangle_\Gamma, \quad \forall \mu_H \in M_H,$$

and balance residual

$$r^{bal} = r - Sr_H.$$

2.  $r_i = D_i^T r^{bal}$  (distribute  $r^{bal}$  to subdomains).

3. Solve local problems for  $\lambda_i \in M_{h,i}$ :

$$S_i \lambda_i = r_i.$$

4.  $\lambda = \sum_{i=1}^n D_i \lambda_i$  (average local solutions).

5. Solve a coarse problem

$$s(\lambda_H, \mu_H) = \langle r, \mu_H \rangle_\Gamma - s(\lambda, \mu_H), \quad \forall \mu_H \in M_H,$$

and balance local solutions

$$M_{bal}^{-1}r = \lambda + \lambda_H.$$

## Condition number estimate

Assume that

$$c\alpha_i \xi^T \xi \leq \xi^T K(x) \xi \leq C\alpha_i \xi^T \xi \quad \forall \xi \in \mathbf{R}^d, \forall x \in \Omega_i.$$

**Lemma**

$$\frac{1}{\alpha_j} s_j(\lambda_i, \lambda_i) \leq \frac{C}{\alpha_j} (1 + \log(H/h))^2 s_i(\lambda_i, \lambda_i).$$

Proof uses:

$$c |\mathcal{I}^{\partial\Omega_i} \mathcal{Q}_{h,i} \mu|_{1/2, \partial\Omega_i}^2 \leq \frac{1}{\alpha_i} s_i(\mu, \mu) \leq C |\mathcal{I}^{\partial\Omega_i} \mathcal{Q}_{h,i} \mu|_{1/2, \partial\Omega_i}^2, \quad \mu \in M_h$$

and

$$\|\mathcal{I}^{\partial\Omega_j} \mathcal{Q}_{h,j} \mu\|_{1/2, \Gamma_{i,j}} \simeq \|\mathcal{I}^{\partial\Omega_i} \mathcal{Q}_{h,i} \mu\|_{1/2, \Gamma_{i,j}}, \quad \mu \in M_h$$

Choose weights on  $\Omega_i$

$$(D_i \lambda_i)(x) = \frac{\alpha_i}{\alpha_i + \alpha_j} \lambda_i(x), \quad x \in \partial\Omega_i \cap \partial\Omega_j.$$

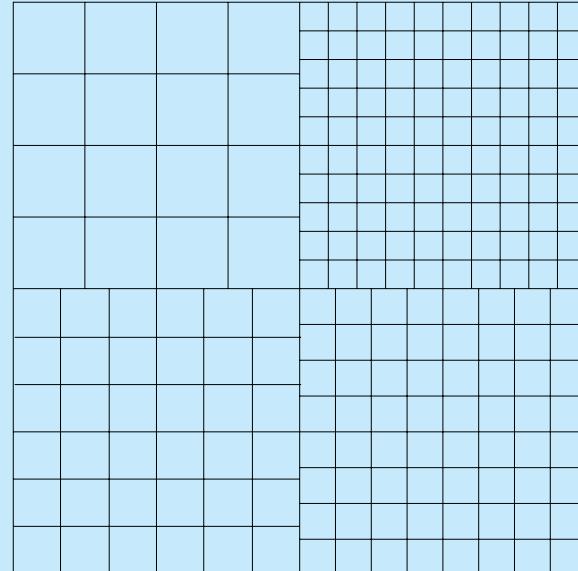
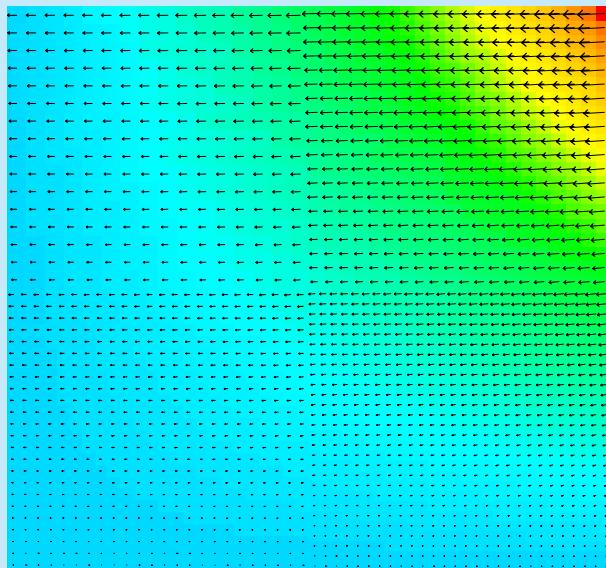
**Theorem**

$$\text{cond}(M_{bal}^{-1} S) \leq C(1 + \log(H/h))^2,$$

where  $C$  does not depend on  $h$ ,  $H$ , and jumps in  $K$ .

## Computational Test 1

$$p(x, y, z) = x^3y^2 + \sin(xy), \quad K = \begin{pmatrix} 10 + 5 \cos(xy) & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

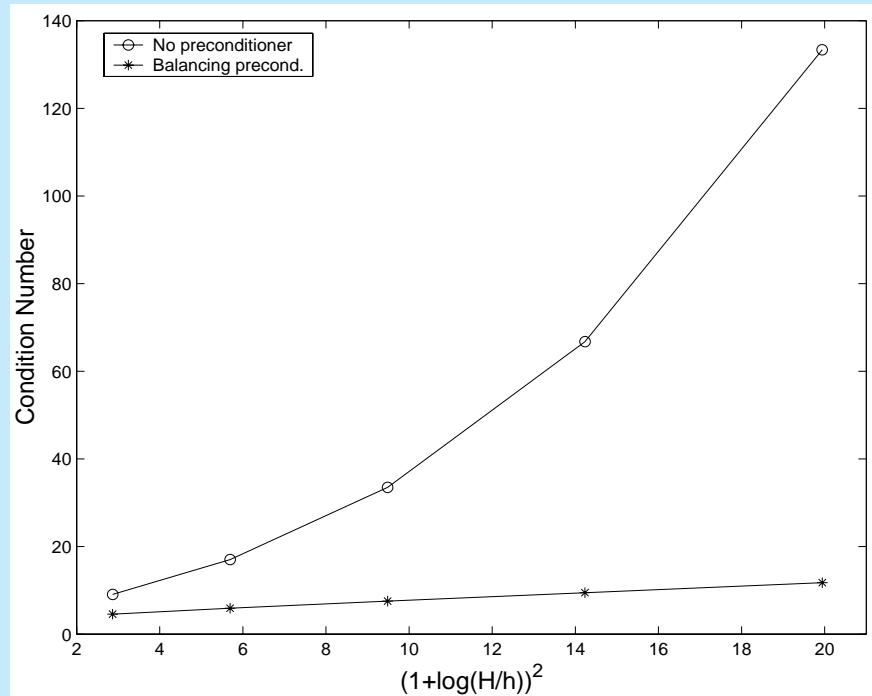
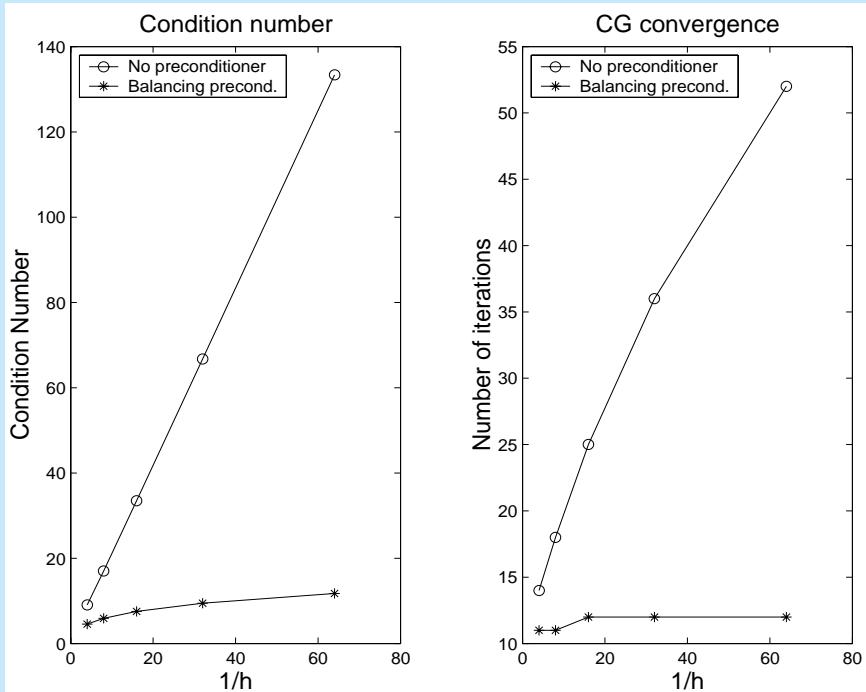


Initial grids

Computed pressure and velocity

## Convergence for Test 1

$1/h$	<i>BalCG</i>		<i>CG</i>	
	<i>cond.</i>	<i>iter.</i>	<i>cond.</i>	<i>iter.</i>
4	4.54520	11	9.04823	14
8	5.90177	11	17.0075	18
16	7.54221	12	33.5087	25
32	9.44828	12	66.7478	36
64	11.7220	12	133.354	52

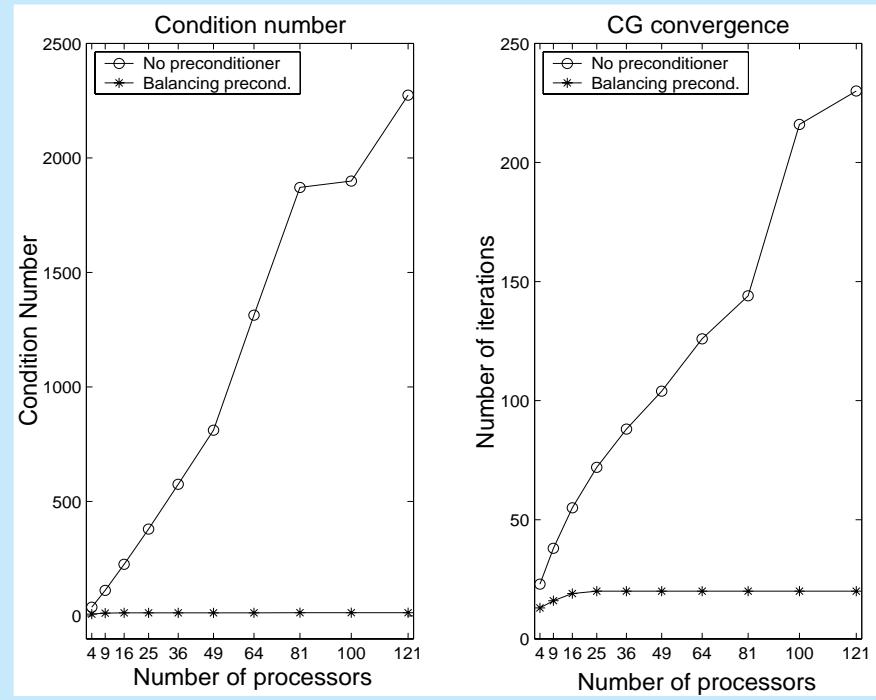


## Computational Test 2

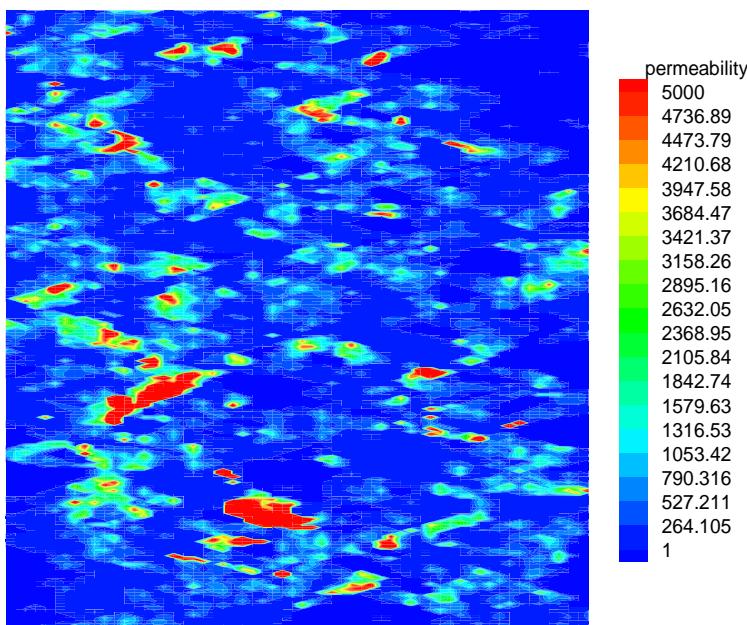
$$p(x, y, z) = x^3y^4 + x^2 + \sin(xy)\cos(y), \quad K = \begin{pmatrix} (x+1)^2 + y^2 & 0 & 0 \\ 0 & (x+1)^2 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

2 x 2 to 11 x 11 subdomains with grids 14 x 18, 12 x 20 in checkboard fashion

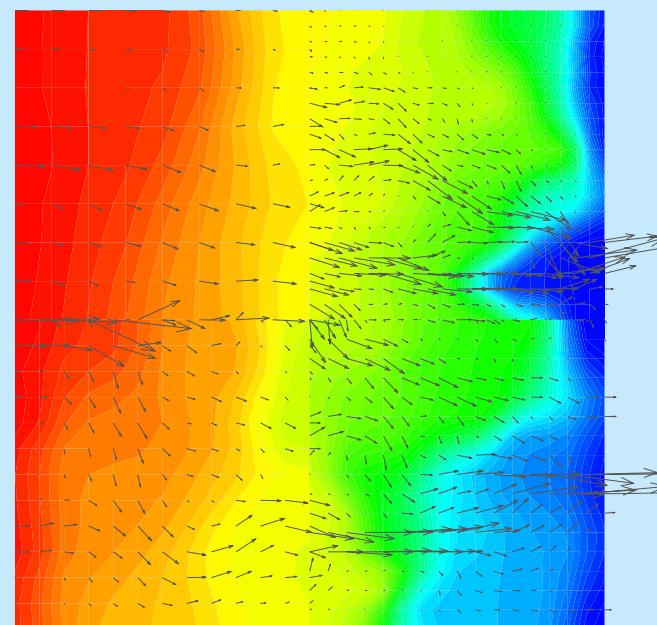
1/h	BalCG		CG	
	cond.	iter.	cond.	iter.
4	8.16598	13	38.2574	23
9	12.2771	16	111.967	38
16	13.2133	19	225.357	55
25	13.4436	20	379.496	72
36	13.6084	20	574.747	88
49	13.6382	20	811.262	104
64	13.7507	20	1313.49	126
81	13.8927	20	1871.29	144
100	13.9882	20	1898.91	216
121	14.0497	20	2274.17	230



## Computational Test 3



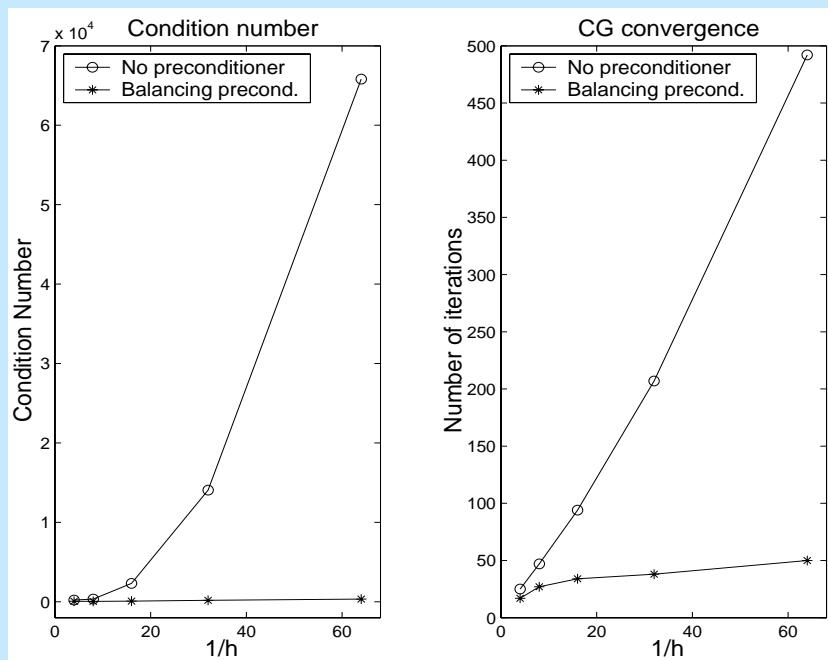
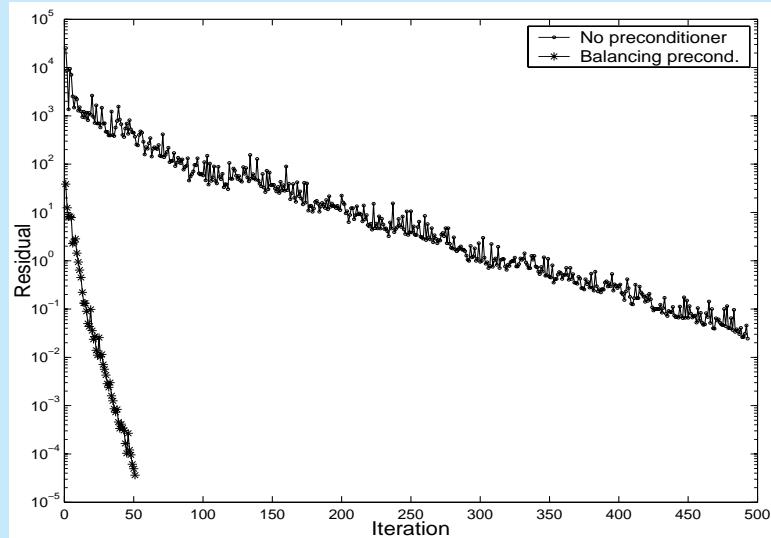
Permeability Tensor  $K$



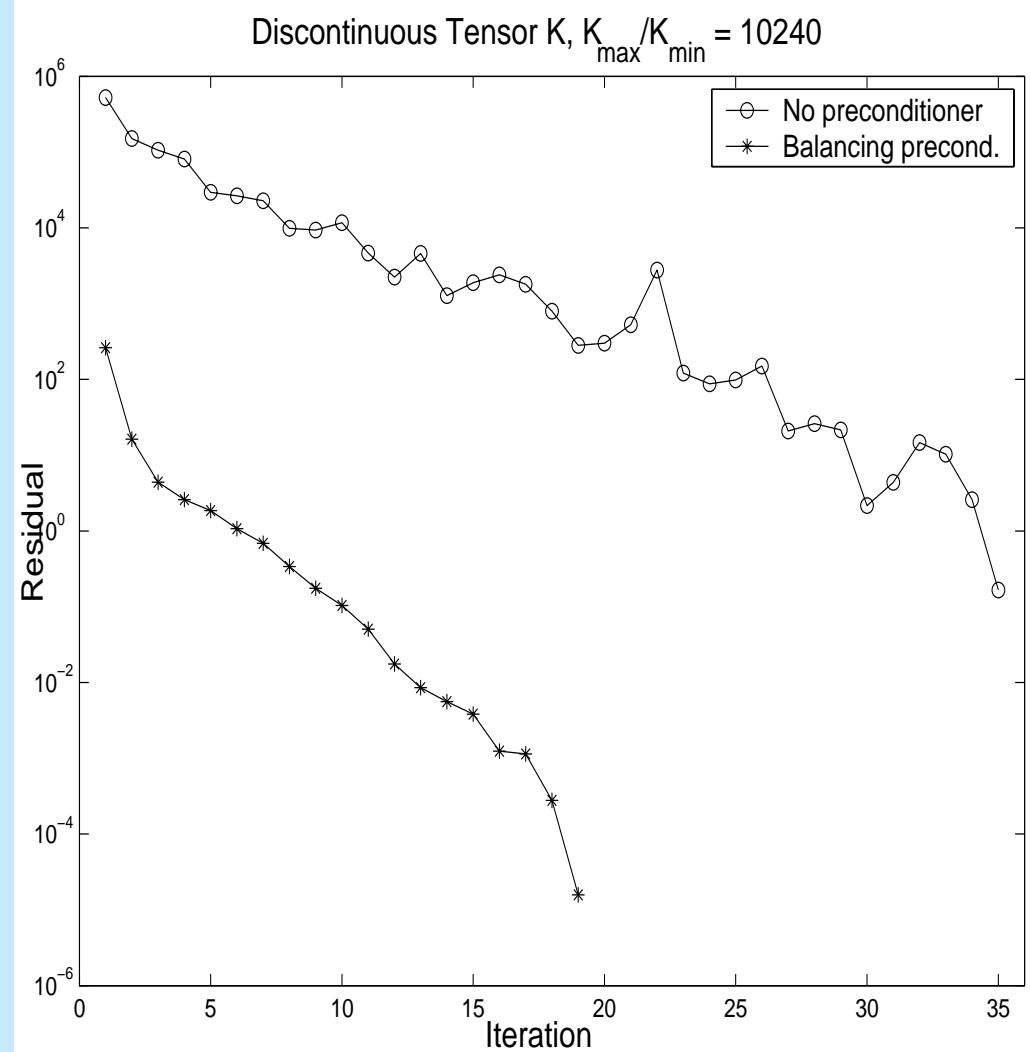
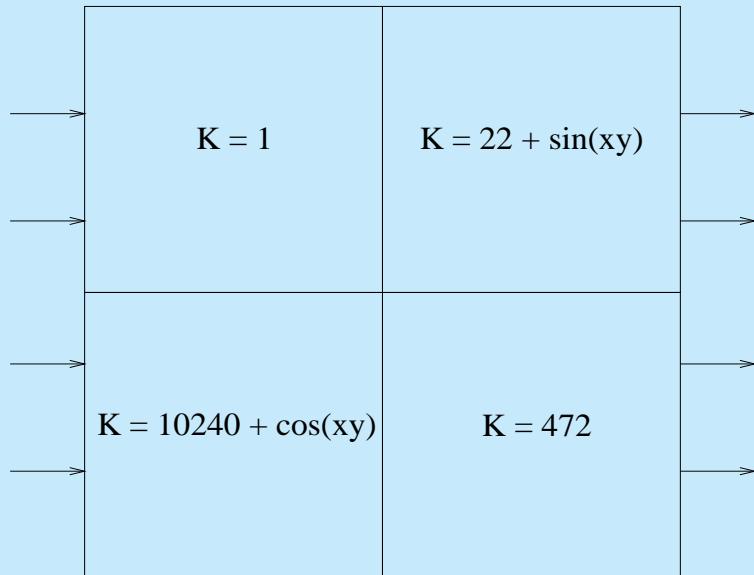
Approximate solution

## Convergence for Test 3

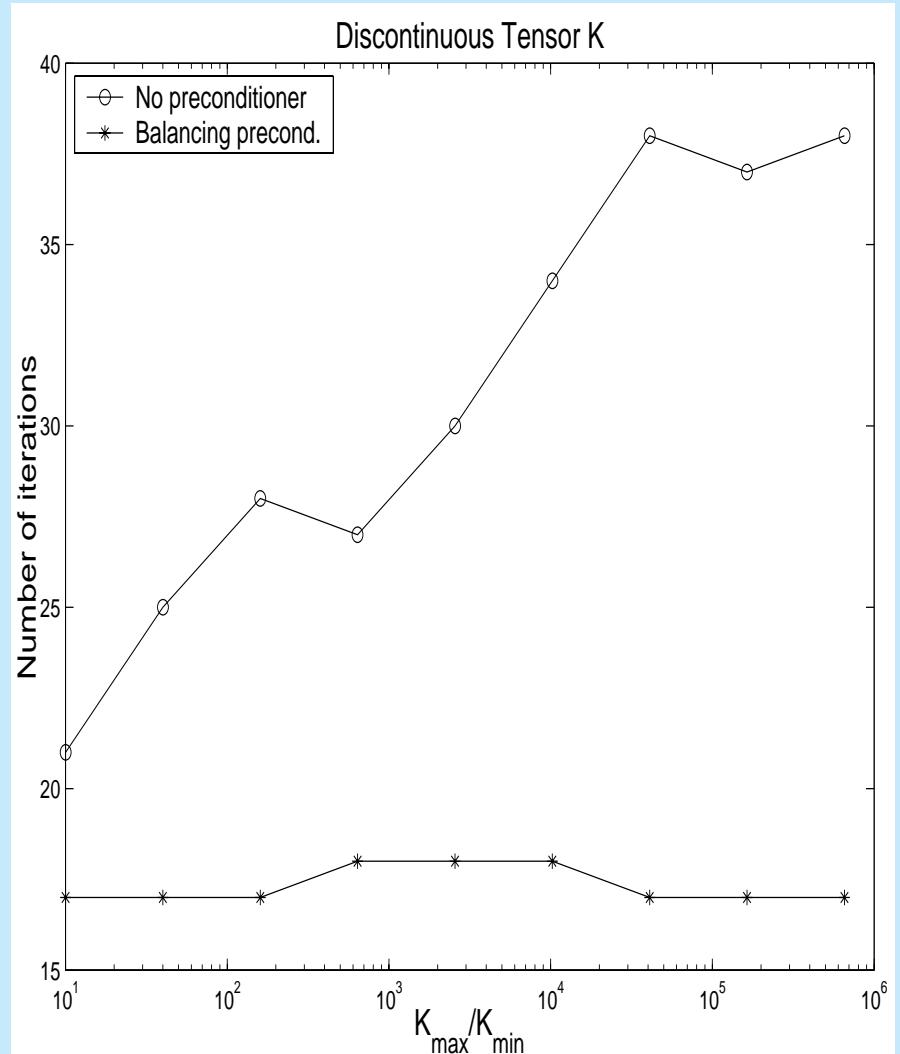
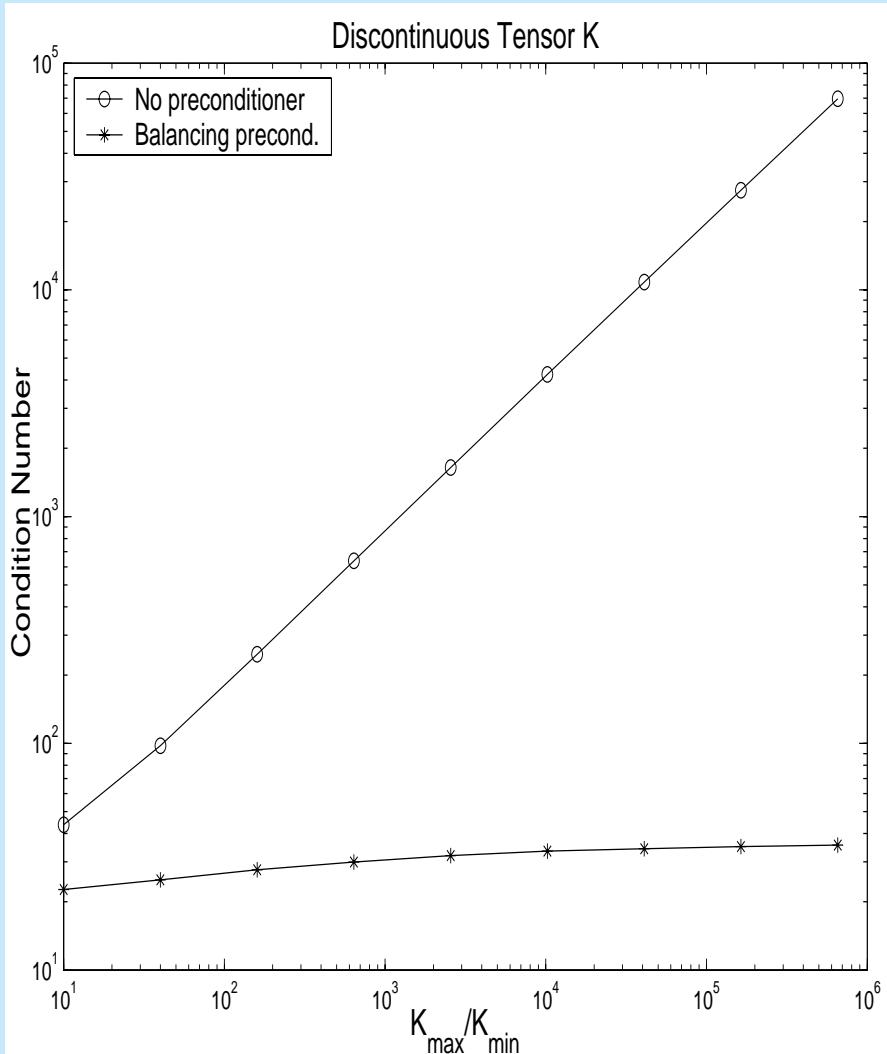
$1/h$	<i>BalCG</i>		<i>CG</i>	
	<i>cond.</i>	<i>iter.</i>	<i>cond.</i>	<i>iter.</i>
4	63.4544	17	222.195	25
8	53.8312	27	338.236	47
16	90.2661	34	2313.30	94
32	192.754	38	14049.8	207
64	336.072	50	65824.2	492



## CG convergence with discontinuous coefficients



## Dependence of CG convergence on jump in coefficients



## Two-phase immiscible flow model

$$\mathbf{U}_\alpha = -\frac{k_\alpha(S_\alpha)K}{\mu_\alpha} \rho_\alpha (\nabla P_\alpha - \rho_\alpha g \nabla D) \quad (\text{Darcy's law})$$

$$\frac{\partial(\phi\rho_\alpha S_\alpha)}{\partial t} + \nabla \cdot \mathbf{U}_\alpha = q_\alpha \quad (\text{conservation of mass})$$

$$S_w + S_n = 1, \quad p_c(S_w) = P_n - P_w$$

$\alpha = w$  (wetting phase),  $n$  (non-wetting phase)

$S_\alpha$  - phase saturation

$\rho_\alpha$  - phase density

$\phi$  - porosity

$K$  - permeability tensor

$k_\alpha(S_\alpha)$  - phase rel. perm.

$g$  - gravitational constant

$N_\alpha = \rho_\alpha S_\alpha$  - concentration

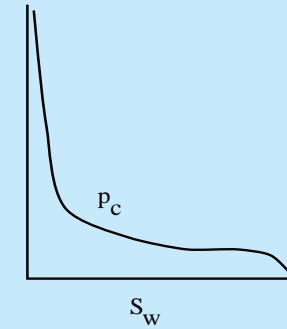
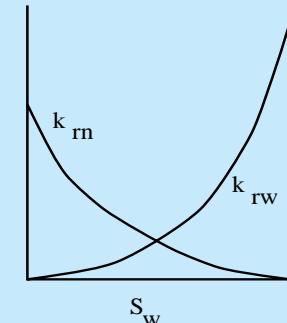
$P_\alpha$  - phase pressure

$q_\alpha$  - source term

$\mathbf{U}_\alpha$  - Darcy velocity

$\mu_\alpha$  - phase viscosity

$D$  - depth



## Multiblock formulation for two-phase immiscible flow model

On each subdomain  $\Omega_i$ :

$$\mathbf{U}_\alpha = -\frac{k_\alpha(S_\alpha)K}{\mu_\alpha} \rho_\alpha (\nabla P_\alpha - \rho_\alpha g \nabla D) \quad (\text{Darcy's law})$$

$$\frac{\partial(\phi\rho_\alpha S_\alpha)}{\partial t} + \nabla \cdot \mathbf{U}_\alpha = q_\alpha \quad (\text{conservation of mass})$$

On each interface  $\Gamma_{ij}$ :

$$P_\alpha|_{\Omega_i} = P_\alpha|_{\Omega_j}, \quad [\mathbf{U}_\alpha \cdot \nu]_{ij} = 0.$$

On each  $\Omega_i$  and  $\Gamma_{ij}$ :

$$S_w + S_n = 1, \quad p_c(S_w) = P_n - P_w.$$

## Parallel domain decomposition - interface formulation

Fully implicit time discretization  $\Rightarrow$   
need to solve a nonlinear system on each time step.

Define interface operator  $B : M \rightarrow M$  for  $\lambda = (P_n^M, P_w^M) \in M$

$$B(\lambda) = ([\mathbf{U}_n^M(\lambda)], [\mathbf{U}_w^M(\lambda)]),$$

where  $\mathbf{U}_\alpha^M(\lambda)$  is the mortar projection of the solution  $\mathbf{U}_\alpha(\lambda) \cdot \nu$  to subdomain problems with Dirichlet boundary data  $\lambda$ .

The original problem is equivalent to solving for  $\lambda \in M$  such that

$$B(\lambda) = 0.$$

## Solution of the interface problem

The interface problem  $B(\lambda) = 0$  is solved by inexact Newton method

$$\lambda_{k+1} = \lambda_k + s_k$$

using forward difference GMRES to approximate the Newton step

$$B'(\lambda_k)s_k = -B(\lambda_k).$$

Each GMRES iteration requires forward difference approximation of  $B'(\lambda)s$ :

$$D_\delta B(\lambda : s) = \frac{B(\lambda + \delta s) - B(\lambda)}{\delta}$$

## Cost of a function evaluation

Algorithm for evaluation of  $B(P_n^M, P_w^M)$ .

1. Project (orthogonally) mortar data onto the subdomain grids

$$P_n^M \xrightarrow{Q_i} \bar{P}_{n,i}, \quad P_w^M \xrightarrow{Q_i} \bar{P}_{w,i}$$

2. Solve in parallel subdomain problems with boundary conditions  $\bar{P}_{n,i}$ ,  $\bar{P}_{w,i}$ .  
Compute  $\mathbf{U}_{n,i}$ ,  $\mathbf{U}_{w,i}$  on each  $\Omega_i$ .

3. Project boundary fluxes back to the mortar space

$$\mathbf{U}_{\alpha,i} \cdot \nu_i \xrightarrow{Q_i^T} \mathbf{U}_{\alpha,i}^M$$

4. Compute flux jump in the mortar space

$$[\mathbf{U}_\alpha^M]_{ij} = \mathbf{U}_{\alpha,i}^M + \mathbf{U}_{\alpha,j}^M$$

## Interface GMRES preconditioner

$D_\delta B$  : Pressure  $\times$  Pressure  $\rightarrow$  Flux  $\times$  Flux

GMRES builds an orthonormal basis for the Krylov space.

Unpreconditioned GMRES step:

$$v_{k+1} = D_\delta B(\lambda : v_k)$$

$v_k \in$  Pressure  $\times$  Pressure;  $v_{k+1} \in$  Flux  $\times$  Flux

$v_k$  and  $v_{k+1}$  live in different physical spaces!

**Preconditioner:**  $M \sim (D_\delta B(\lambda))^{-1}$

$M$ : Flux  $\times$  Flux  $\rightarrow$  Pressure  $\times$  Pressure

Preconditioned GMRES step:

$$v_{k+1} = M D_\delta B(\lambda : v_k)$$

$v_{k+1} \in$  Pressure  $\times$  Pressure

$v_k$  and  $v_{k+1}$  live in the same physical space.

The preconditioned GMRES iterates in the Pressure  $\times$  Pressure space.

## Neumann-Neumann preconditioner for $D_\delta B$

$$D_\delta B = \sum_i D_\delta B_i$$

$D_\delta B$  : Dirichlet to Neumann operator (Poincare-Steklov)

Preconditioner:

$$M = \sum_i \widehat{D_\delta B}_i^{-1},$$

where  $\widehat{D_\delta B}_i^{-1}$  is an approximation to  $D_\delta B_i^{-1}$ .

$M$  : Neumann to Dirichlet operator

## Approximation to the Pressure → Flux interface operator

On  $\Gamma_{12}$ ,  $B = B_1 + B_2$ ,

$B_i$  : Pressure → Flux is based on Darcy's Law

$$\mathbf{U} \cdot \nu_i = -(K \nabla P) \cdot \nu_i$$

Approximate  $B_i$  by a finite difference on a single-cell layer along the interface:

$$\hat{B}_i(\bar{P}) = K_i \frac{P_i - \bar{P}}{h_i/2}$$

Leading to an approximation to  $D_\delta B(\bar{P} : s)$

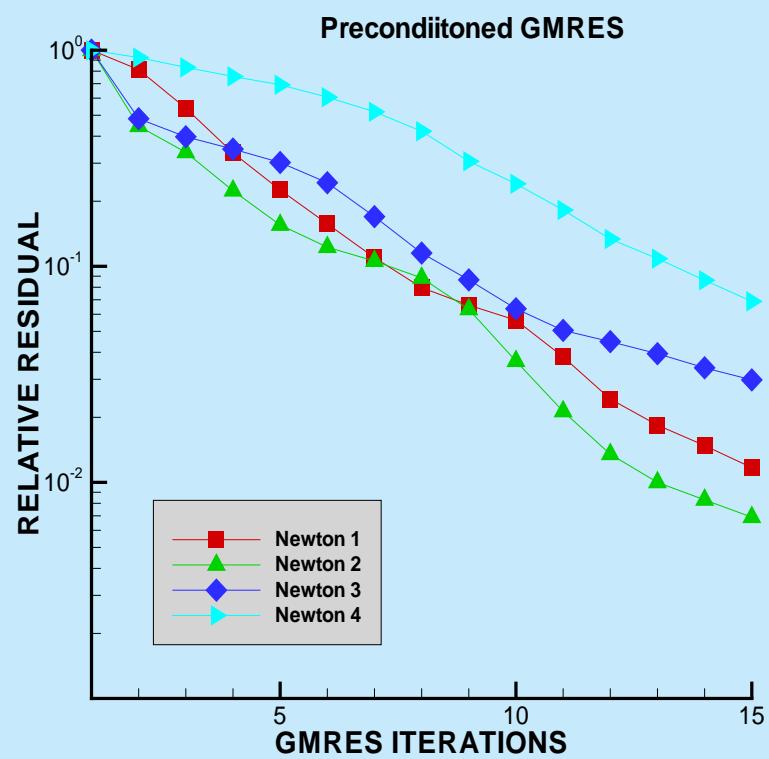
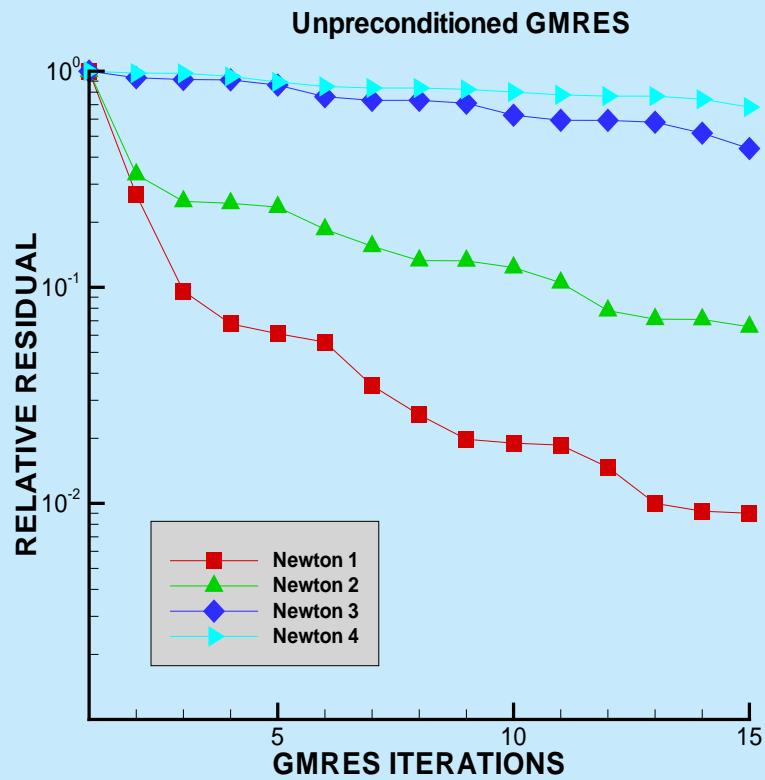
$$\widehat{D_\delta B}(\bar{P} : s) = \frac{\hat{B}(\bar{P} + \delta s) - \hat{B}(\bar{P})}{\delta} = -2\left(\frac{K_1}{h_1} + \frac{K_2}{h_2}\right)s$$

Preconditioner: given flux  $v$

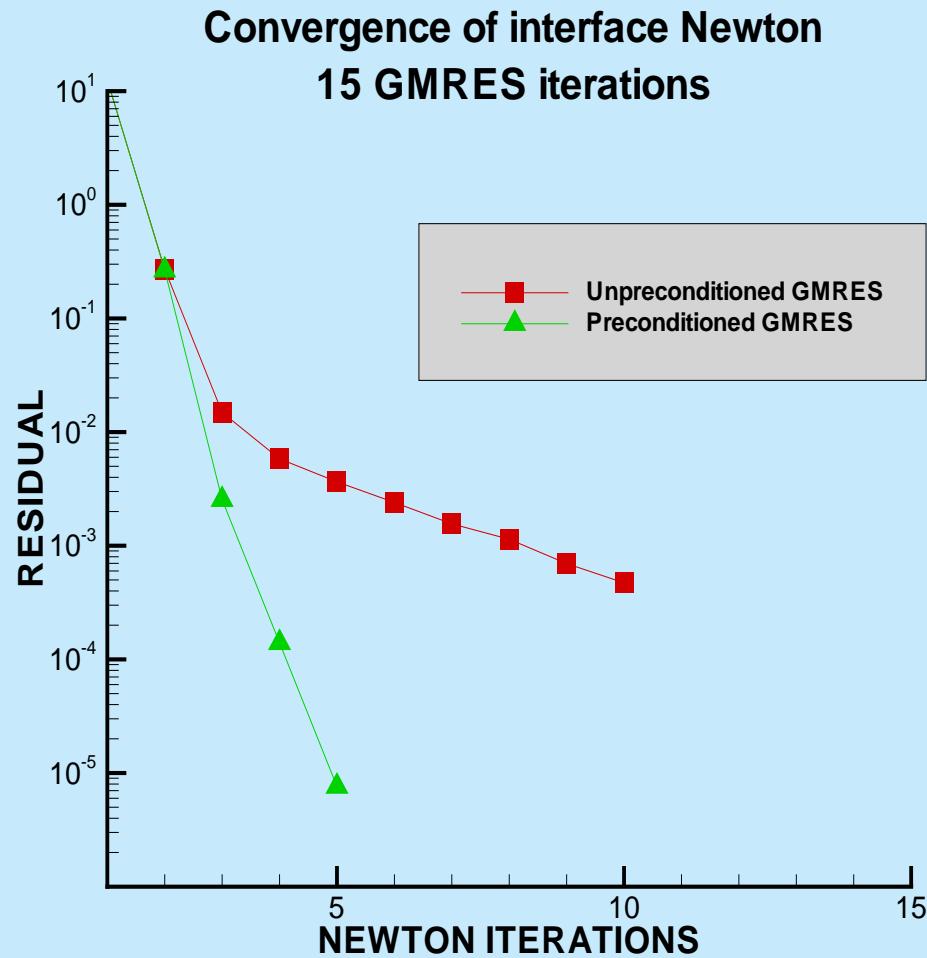
$$Mv = -\frac{1}{2\left(\frac{K_1}{h_1} + \frac{K_2}{h_2}\right)}v$$

## Interface GMRES convergence

Using the CSM simulator IPARS.



## Effect of preconditioner on interface Newton convergence



## **Conclusions**

- Multiblock formulation on non-matching grids
- Domain decomposition - interface problem
- Balancing preconditioner for single-phase flow
  - Consistent Neumann solves
  - Coarse solve provides global exchange of information
  - Quasi-optimal condition number independent of the jump in coefficients
- Newton-GMRES with Neumann-Neumann preconditioner for multiphase flow

## **Current work**

- Balancing for multiphase flow
- Direct linearization of the global system
- IMPEs-type formulation